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The Conformation of Non-Aromatic Ring Compounds. VIII.* The Crystal Structure of *cis*-2,3-Dichloro-1,4-dioxane at -140 °C

BY C. ALTONA AND C. ROMERS

Laboratory of Organic Chemistry, University of Leiden, The Netherlands

(Received 28 January 1963)

cis-2,3-Dichloro-1,4-dioxane crystallizes in the orthorhombic system, space group $P2_12_12_1$, Z = 4, and

a = 4.463, b = 10.67, c = 13.13 Å at -140 °C.

The structure was refined in three projections by means of data obtained at about -140 °C with molybdenum radiation.

The dioxane ring has the chair form. The two carbon-chlorine bonds are unequal in length: C(1)-Cl(1) (axial) = 1.819 ± 0.009 Å, C(2)-Cl(2) (equatorial) = 1.781 ± 0.007 Å.

Introduction

It seemed worthwhile to determine the structure of cis-2,3-dichloro-1,4-dioxane by X-ray methods, since this compound is the only known dihalogeno-1,4-dioxane in which an equatorial carbon-halogen bond was expected to be present (Altona, Romers & Havinga, 1959).

Experimental

cis-2,3-Dichloro-1,4-dioxane was prepared by chlorination of a refluxing mixture of 1,4-dioxane and carbon tetrachloride (Summerbell & Lunk, 1957). The compound crystallizes from carbon tetrachloride as colourless needles, elongated along the *a* axis. The crystals are soft and tend to disintegrate when even slight pressure is applied in directions perpendicular to the needle axis. After several trials some good crystals were obtained that could be cut along [010] and [001]. Their cross sections were roughly hexagonal with faces {001} and {011} about equally developed. A central capillary hole with a radius of approximately 0.02 mm ran along the full length of all these needles.

The dimensions of the unit cell were determined from zero-layer Weissenberg photographs taken with unfiltered Cu $K\alpha$ radiation ($\lambda = 1.5418$ Å) about [100],

^{*} Part VII: Altona, Knobler & Romers (1963); for earlier parts see van Dort & Havinga (1963).

[010], and [001] at -140 °C. The powder diffraction lines of aluminum ($\alpha = 4.0189$ Å at 20 °C) were superposed on the films.

Crystal data

cis-2,3-Dichloro-1,4-dioxane, C₄H₆O₂Cl₂, m.p. 53·5– 54 °C. Orthorhombic, $a = 4.463 \pm 0.01$, $b = 10.72 \pm 0.02$, c = 13.13 Å, measured at -140 °C.

 D_x from data obtained at 20 °C (Z=4), 1.63 g.cm⁻³; D_m at 20 °C, 1.59 g.cm⁻³. Absorption coefficient for Mo K α rays (λ =0.7107 Å) μ =9.4 cm⁻¹. Space group, from absent spectra, $P2_{12}2_{1}$.†

Table 1. Technical data concerning the intensity records

| | Expo- | \mathbf{Number} | | | |
|-----------------|--------------------|-------------------|----------------------------------|---------|------------------|
| | sure | \mathbf{of} | | Temper- | Cross |
| | $_{\mathrm{time}}$ | observed | $(\sin\theta/\lambda)_{\rm max}$ | ature | section |
| \mathbf{Zone} | (hr) | reflexions | (Å-1) | (°C) | (mm) |
| [100] | 77 | 247 | 0.96 | -130 | 0.4×0.4 |
| [010] | 94 | 140 | 1.10 | -142 | 0.4×0.5 |
| [001] | 71 | 90 | 1.03 | -142 | 0.6×0.6 |

Intensity records were obtained on zero-layer Weissenberg exposures about [100], [010], and [001], taken with filtered Mo $K\alpha$ radiation at low temperatures (Table 1). The multiple-film technique was used with thin copper foils inserted between films. The crystal specimens were sealed in Lindemann glass capillaries. The visually estimated intensities were corrected for Lorentz and polarization factors by the usual methods. Absorption effects were neglected since $\mu R = 0.38-0.56$ cm⁻¹ for Mo $K\alpha$ radiation. Preliminary scale and temperature factors changed less than 3% in later stages of the analysis; the absolute scale was established by correlation of the observed with the calculated structure factors. Additional technical data are collected in Table 1.

Determination and refinement of the structure

The Patterson functions P(vw) and P(uw) were calculated from preliminary Mo $K\alpha$ data obtained at room temperature. From the Patterson maps the positions of the chlorine atoms could be determined without difficulty once it was recognized that the two atoms almost completely overlap in the [010] projection (Fig. 1(b)). An electron density projection $\rho(yz)$, calculated with 65% of the structure factors whose signs were based upon the chlorine contributions, clearly showed the carbon and oxygen atoms. This method worked less well for the projection along [010], because the chlorine atoms contribute only to about half of the structure factors F(h0l). Since the main features of the molecule were now known, the x parameters of the atoms in the dioxane ring could be taken from a model. An electron density map



Fig. 1. Electron-density projections of cis-2,3-dichloro-1,4-dioxane. (a) Projection along [100]; contour lines around the carbon and oxygen atoms are drawn at 0, 2, 4, ... e.Å⁻², those around the chlorine atoms at 0, 2, 5, 10, ... e.Å⁻².
(b) Projection along [010]; contour lines around C: 5, 7.5, 10, ... e.Å⁻²; around Cl: 5, 10, 15, ... e.Å⁻².
(c) Projection along [001]; contour lines around C and O: 5, 8, 10, ... e.Å⁻²; around Cl: 5, 8, 12, ... e.Å⁻².

 $\varrho(xz)$ showed that this model was essentially correct; however, serious overlap of atomic peaks in the [010] and [001] projections soon prevented further progress.

In view of our favourable experiences with low

[†] In a preliminary paper (Altona, Romers & Havinga, 1959) the space group was erroneously reported to be $P2_1/n$.

Table 2. Observed and calculated structure factors

In each column the running index is followed by $10 F_o$ and $10 F_c$, in that order. The non-observed reflexions, indicated by x, are inserted at the minimum observable value.

| 0 | K L | 04Ì | 0 9 l | 0 14 L | 1 0 | 5 o l | 1 k 0 |
|---|--|---|--|---|--|--|---|
| 0 2 624 4 571 6 72 8 10 10 371 12 224 14 99 16 23 20 84 22 90 0 | 0 8 6 13 7 7 140 3 98 0 -395 6 237 0 105 0 -240 9 33 6 78 0 -80 -80 -155 -551 -551 -251 -251 -251 -255 -256 | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $ \begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ |
| 1 X 2 35(3 5) 4 39(5 197) 6 107 7 102 8 217 9 74 10 216 11 37 12 34(13 25 14 25 15 6 32 16 32 17 55 | $\begin{array}{c} 55 \\ -397 \\ 8 \\ -50 \\ 5 \\ 415 \\ 7 \\ 203 \\ 111 \\ 9 \\ -296 \\ 130 \\ 3 \\ 111 \\ 9 \\ -296 \\ 237 \\ 7 \\ 38 \\ 7 \\ 12 \\ 9x \\ -17 \\ 2x \\ 15 \\ 35 \\ 15 \\ 5 \\ 3 \\ 15 \\ 5 \\ 3 \\ 15 \\ 5 \\ 15 \\ 5 \\ 15 \\ 5 \\ 15 \\ 5 \\ 15 \\ 5 \\ $ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ |
| 10 95 19 62 0 488 1 12 2 518 3 140 5 259 7 144 8 25 9 78 10 144 11 27 | 2 2 3 -445 -54 -143 -143 -143 -147 -147 -261 -137 -28 -33 -149 -24 | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | 1 54 56 2 360 351 3 227 231 4 467 438 5 66 71 6 385 -396 7 28z -13 8 277 -288 9 29z - 32 10 252 262 11 62 - 59 12 37 34 17 83 - 78 4 ± 0 |
| 12 78 13 90 14 33 15 62 16 74 17 78 19 38 20 38 21 66 0 1 283 2 756 | -97 -97 -97 -74 -74 -74 -72 -723 | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $ \begin{array}{cccccccccccccccccccccccccccccccccccc$ |
| 3 510 4 629 5 111 6 353 7 131 8 407 9 201 10 37 11 53 12 152 13 181 14 164 15 33 17 123 16 53 17 123 18 152 19 38 20 62 | $\begin{array}{c} -507\\ 613\\ -110\\ 366\\ 135\\ -423\\ -224\\ 52\\ -224\\ 52\\ -751\\ 173\\ 203\\ -169\\ -36\\ -118\\ 141\\ 141\\ x \\ 15\\ -50\end{array}$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{c} c \ 19 \ 1 \\ 1 \ 42x \ 15 \\ 2 \ 42 \ -44 \\ 0 \ 20 \ 1 \\ 0 \ 49 \ -52 \\ 1 \ 42x \ -25 \\ 2 \ 43 \ 37 \\ 3 \ 42x \ 0 \\ 49 \ 6 \ 42 \ -39 \end{array}$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{c} 5 & 32x - 42 \\ 32x - 5 \\ 7 & 32x - 8 \\ 8 & 32 - 31 \\ 9 & 32x - 3 \\ 10 & 78 & 50 \\ 9 & 0 \\ 8 & 39 & 50 \\ 9 & 35 & -28 \\ \hline \\ H & K & 0 \end{array}$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ |
| 21 38 22 38 23 38 24 49 0 0 908 1 501 | $4 \begin{bmatrix} -901 \\ -972 \end{bmatrix}$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$ | E O L O O L 2 600 -616 4 546 -558 6 753 -764 | $\begin{array}{c} 4 & 0 \\ 0 & 324 & 331 \\ 1 & 226 & 227 \\ 2 & 179 & 162 \\ 3 & 21x - 20 \\ 4 & 308 & -297 \\ 5 & 109 & 132 \\ 6 & 214 & -206 \\ 7 & 121 & -137 \end{array}$ | 2 488 451 4 906 -909 6 62 - 74 8 70 76 10 277 316 12 31x - 16 14 294 -316 | 6 k 0 0 132 -130 2 79 75 3 30 - 38 4 145 161 7 74 73 |
| 2 140 3 86 4 70 5 452 6 509 7 131 8 82 9 90 10 232 11 201 | -160 -93 48 444 -500 -132 -89 -100 238 222 | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $ \begin{array}{ccccccccccccccccccccccccccccccccccc$ | 8 23 - 47 9 82 - 76 10 246 240 11 24x 1 12 168 186 13 86 96 14 125 -137 15 43 48 16 133 -138 17 78 - 86 | $\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$ | 7×0 $1 38 -28$ $2 41 25$ $3 62 55$ $4 45 38$ $6 62 -48$ $10 66 61$ |
| 12 107 13 58 14 70 15 90 16 164 | -120 - 43 - 75 - 94 - 174 | 1 185 -180 2 26x 16 3 164 165 4 27 26 5 74 - 86 | 0 242 261 1 86 - 86 2 34x - 3 3 34x - 29 | 1 0 L 1 717 -711 | 18 30x 13 19 30x 6 20 47 39 21 33x - 4 22 47 59 | 5 95 - 97 6 496 529 7 24x 7 8 136 136 5 66 50 | 8 k 0 0 66 53 3 66 54 7 83 - 70 |



Fig. 2. Difference Fourier syntheses. (a) The final $F_o - F_c$ projection along [100], showing the principal axes of the thermal vibration ellipses of the chlorine atoms. Contour lines are drawn at 0, 0.4, 0.8, ... e.Å⁻², negative contours are broken. (b) Region around chlorine atom Cl(1) before a correction was applied (see text) for the thermal anisotropy. Contours are drawn at approximately 0, 0.3, 0.6, ... e.Å⁻²; negative contours are broken.

temperature data in the analysis of the structure of the trans-2,5-isomer (Altona, Knobler & Romers, 1963), we decided to continue the refinement by using a new set of Mo $K\alpha$ reflexions collected at -140 °C. Calculation of two cycles of structure factors and Fourier syntheses (Fig. 1) yielded the following reliability factors R and temperature constants B:

| F(0kl), | R = 14.9%, | $B = 2 \cdot 32 \text{ Å}^2;$ |
|---------|------------|-------------------------------|
| F(h0l), | R = 14.7%, | $B = 1.77 \text{ Å}^2;$ |
| F(hk0), | R = 13.2%, | $B = 1.98 \text{ Å}^2$. |

The structure was further refined in four successive cycles of structure factors and difference syntheses with contributions of the heavy atoms subtracted. Atomic positions and individual isotropic temperature factors were adjusted. Except in cases of severe overlap, each projection was treated independently; this was done partly in order to obtain a check upon the accuracy of the positional parameters and partly because systematic discrepancies were noted between the observed high-angle F(00l) values derived from the (0kl) and from the (h0l) films respectively (Table 2). These discrepancies may be due to photographic errors or to a fairly large temperature difference between the two crystals. Such a difference could occur if the a-axis crystal had not been placed in the centre of the cold gas stream during the exposure.

The thermal vibrations of the chlorine atom were markedly anisotropic (Fig. 2(b)). A correction for this effect was applied by substituting for each chlorine atom two half-atoms placed 0.08-0.12 Å apart in the direction of maximum vibration (Kartha & Ahmed, 1960; Vos & Smits, 1961). The final separations were about 15% larger than those calculated initially from the formulae proposed by Kartha & Ahmed (1960). Although this method is approximative, satisfactory results were obtained; in the final difference maps no significant electron density details remained in the vicinity of the chlorine atoms (Fig. 2(a)).

In an attempt to find reliable hydrogen atom positions, difference maps were calculated in which only low-angle reflexions were used $(\sin \theta / \lambda < 0.55 \text{ Å}^{-1})$. The hydrogen maxima showed up clearly with peak heights of 0.6-0.8 e. Å⁻². Because of overlap, however, some hydrogen parameters could not be determined. We therefore preferred to insert these atoms into the molecular model by accepting a C-H distance of 1.03 Å and tetrahedral angles.

The final R indices, excluding non-observed reflexions, were:

R(0kl) = 6.3%, R(h0l) = 7.3%, R(hk0) = 7.1%.

The atomic scattering factors for carbon and oxygen were those calculated by Berghuis, Haanappel, Potters, Loopstra, MacGillavry & Veenendaal (1955), for chlorine those by Tomiie & Stam (1958), and for hydrogen those by McWeeny (1951). The difference Fourier syntheses (at intervals of 1/128 of the cell edges) and the corresponding structure factors were calculated on a ZEBRA computer with programs devised by Smits (1961). A list of the observed and calculated structure factors is given in Table 2. In order to retain a centre of symmetry in each projection, the three-dimensional origin was shifted according to the following scheme: [100], $(0, -\frac{1}{4}, 0)$; [010], $(0, 0, \frac{1}{4})$; [001], $(\frac{1}{4}, 0, 0)$.

The agreement between the positional parameters derived from each projection was satisfactory. Eleven of the fourteen parameters that were determined twice showed differences less than the sum of the respective standard errors. None of the remaining differences was possibly significant. Since the projection along [001] suffered from overlap effects (Fig. 1(c)), the x parameters were taken from the [010] projection and the y values from the [100] projection. The z parameters of O(2), O(1), C(1), $\tilde{C}(2)$ and C(4) are weighted means. The final atomic positions are shown in Table 3. Their standard errors were estimated by Cruickshank's (1949) method. Since only two-dimensional methods were employed, these standard errors might be too low and should perhaps be increased by about 50%.

Table 3. Final atomic coordinates (in fractions of the cell edges) and standard deviations

| ł | Å | × | 10 |)-3) |
|---|---|---|----|------|
| ۰ | A | ~ | T | , , |

| Atom | x | y | z |
|-------|--------------|-------------|-------------|
| Cl(1) | -0.0129 (3) | 0.1206(2) | 0.1515(2) |
| Cl(2) | -0.0161 (3) | -0.1785(2) | 0.1580(2) |
| O(1) | -0.3000 (9) | 0.0945(6) | -0.0252 (5) |
| O(2) | 0.1555(9) | -0.0916 (5) | -0.0140 (5) |
| C(1) | -0.2470 (11) | 0.0302 (8) | 0.0653 (7) |
| C(2) | -0.1140 (10) | -0.0985 (7) | 0.0439 (6) |
| C(3) | 0.0900(11) | -0.0279(7) | -0.1109 (7) |
| C(4) | -0.0300 (13) | 0.1012 (8) | -0.0886 (7) |
| H(1) | -0.441 | 0.005 | 0.100 |
| H(2) | -0.267 | -0.120 | 0.001 |
| H(3) | 0.264 | -0.015 | -0.159 |
| H(4) | -0.038 | -0.018 | -0.159 |
| H(5) | -0.096 | 0.141 | -0.160 |
| H(6) | 0.105 | 0.160 | -0.050 |

The central curvatures and the peak electron densities are shown in Table 4. These values are of the same order of magnitude as those found in the low temperature projections of *trans*-2,5-dichloro-dioxane. The individual isotropic temperature factors, as well as the values of B_{\min} and B_{\max} for each chlorine atom, are also given in Table 4. The directions of the vibrational axes, projected along [100],

Table 4. Isotropic temperature factors B (Å²), peak electron densities ρ_0 (e.Å⁻²) and average central curvatures (e.Å⁻⁴)

Two values are shown for each chlorine atom, the one between brackets pertains to the component in the direction of maximum vibration

| | [100] | | | [010] | | | [001] |
|-------|--------------|--------------|---------------------------------|--------|--------------|---------------------------------|--------|
| | B | Q0 | $\delta^2 \varrho / \delta r^2$ | В | | $\delta^2 \varrho / \delta r^2$ | B |
| Cl(1) | 1.65 | 3 8·7 | 675 | 1.65 | * | * | 1.18 |
| | (2.50) | | (560) | (3.05) | | | (2.08) |
| Cl(2) | 1.68 | 38.9 | 625 | 1.65 | * | * | 1.48 |
| | (2.44) | | (485) | (3.05) | | | (2.48) |
| O(1) | 2.29 | 15.5 | 200 | 1.17 | 20.8 | 305 | 1.38 |
| O(2) | $2 \cdot 01$ | 15.7 | 215 | 1.00 | $22 \cdot 6$ | 320 | 1.50 |
| C(1) | 1.88 | 10.6 | 140 | 1.31 | 15.6 | 255 | 1.22 |
| C(2) | 1.92 | 12.1 | 185 | 0.90 | 17.0 | 300 | 1.54 |
| C(3) | 1.95 | 11.2 | 160 | 1.12 | 15.3 | 250 | 1.54 |
| C(4) | $2 \cdot 17$ | 11.0 | 145 | 1.25 | 14.3 | 210 | 1.62 |
| H | 2.00 | | | 1.00 | | | 1.30 |

* The chlorine atoms suffer from overlap in the projection along [010].

are shown in Fig. 2. In the [010] and [001] projections the maximum vibration vector almost coincided with the x axis.

Conformation of the molecule

The dioxane ring has the chair form with one chlorine atom (Cl1) in the axial and the other (Cl2) in the equatorial position (Fig. 3). This conformation was predicted earlier from dipole moment data (Altona, Romers & Havinga, 1959). The bond distances and angles are shown in Table 5 (see also Fig. 3). The important features can be summarized as follows:

| Table 5. Intramolecular | distances, | angles | and |
|-------------------------|---------------|------------------|-----|
| estimated stands | ard deviation | ons [¯] | |

| Bond | Distance | σ |
|----------------------|---------------|----------------------|
| C(1) - Cl(1) | 1·819 Å | 9×10^{-3} Å |
| C(2)-Cl(2) | 1.781 | 7 |
| C(1) - C(2) | 1.528 | 11 |
| C(3) - C(4) | 1.513 | 11 |
| C(1) - O(1) | 1.394 | 9 |
| C(2) - O(2) | 1.425 | 12 |
| C(3) - O(2) | 1.466 | 9 |
| C(4)-O(1) | 1.473 | 12 |
| Non-bonded | | |
| $CI(1) \cdots O(1)$ | 2.664 | 7 |
| $Cl(1) \cdots Cl(2)$ | $3 \cdot 207$ | 4 |
| Bonds | Angle | σ |
| O(1) - C(1) - C(2) | 110·8° | 0.7° |
| C(1) - C(2) - O(2) | 112.3 | 0.7 |
| O(2) - C(3) - C(4) | 109.1 | 0.7 |
| C(3)-C(4)-O(1) | 110.8 | 0.7 |
| C(1)-O(1)-C(4) | 111.6 | 0.6 |
| C(2) - O(2) - C(3) | 108.5 | 0.6 |
| Cl(1)-C(1)-O(1) | 111.3 | 0.5 |
| Cl(1)-C(1)-C(2) | 111.8 | 0.6 |
| Cl(2)-C(2)-C(1) | 112.0 | 0.6 |
| Cl(2)-C(2)-O(2) | 105.5 | 0.6 |



Fig. 3. Atomic distances and bond angles in cis-2,3-dichloro-1,4-dioxane.

(1) The axial carbon-chlorine bond distance $(1\cdot819\pm0\cdot009 \text{ Å})$ is significantly larger $(p=0\cdot2\%)$; Cruickshank & Robertson, 1953) than the equatorial one $(1\cdot781\pm0\cdot007 \text{ Å})$.

(2) The bond C(1)–O(1) $(1.394 \pm 0.009 \text{ Å})$, adjacent to the axial C(1)–Cl(1) bond, is perhaps (p=1-5%) shorter than the carbon-oxygen bond in dimethyl ether $(1.416 \pm 0.003 \text{ Å})$, Kimura & Kubo, 1959). The bond distance of C(2)–O(2), adjacent to the equatorial C(2)–Cl(2) bond, appears to be normal $(1.425 \pm 0.012 \text{ Å})$.

(3) The bond distances of C(4)–O(1) $(1.466 \pm 0.013 \text{ Å})$ and C(3)–O(2) $(1.473 \pm 0.010 \text{ Å})$ are relatively large and deviate significantly (p=0.1%) from the C–O bond distance in dimethyl ether quoted above.

(4) The 'equatorial half' of the molecule, consisting of the atoms Cl(2), C(2), O(2) and C(3), does not deviate significantly from planarity. The equation of the plane, obtained by a least-squares procedure, is:

0.2758X + 1.1035Y + 0.5421Z + 1 = 0.

All four atoms are at distances less than 0.007 Å from this plane. The angles between the bonds lying in the plane, Cl(2)-C(2)-O(2) ($105.5\pm0.6^{\circ}$) and C(2)-O(2)-C(3) ($108.5\pm0.6^{\circ}$) are smaller than the corresponding angles in the 'axial half'.

Not too much stress should be laid on relatively small differences in bond lengths in view of the uncertainties in the estimation of two-dimensional standard errors (see preceding section).

One might expect some steric strain to be present in this molecule, owing to the close approach of the two chlorine atoms (3.207 Å). However, it is remarkable that, compared with the conventional chair model (tetrahedral angles), the vectorial angle between the C-Cl bonds is narrowed from 70.5 to 66°, and the dihedral angle is only 50.5 instead of 60°.

The bond distances and angles in the 'axial half' of the molecule resemble those occurring in the *trans*-2,5isomer, in which compound a long axial C-Cl bond (1.845 Å) is present adjacent to a short C-O bond (1.388 Å; Altona, Knobler & Romers, 1963). Similar effects occur in *trans*-2,3-dichloro-1,4-dioxane (Altona & Romers, 1963).

Molecular environment

All intermolecular atomic distances <4.4 Å have been calculated on the X₁ computer, University of Leiden, using a program (A 200) written in ALGOL 60 by Mr H. L. Jonkers. A selection of these distances is listed in Table 6 (see also Fig. 4).

| Tabl | е6. | Some | intermolecu | ular o | $distances^*$ | and | angles |
|------|-----|------|-------------|--------|---------------|-----|--------|
|------|-----|------|-------------|--------|---------------|-----|--------|

| Atoms | | 3 | Distar | nce |
|-----------------------|-------------------------------------|-----|--------|-------|
| $O(2) \cdots O(1)$ | $[0]_{x}$ | | 3.14 | Å |
| $O(2) \cdots C(1)$ | $\begin{bmatrix} 0 \end{bmatrix}_x$ | | 3.14 | |
| $O(2) \cdots H(1)$ | $\begin{bmatrix} 0 \end{bmatrix}_x$ | | 2.55 | |
| $O(2) \cdots H(2)$ | $\begin{bmatrix} 0 \end{bmatrix}_x$ | | 2.65 | |
| $C(3) \cdots O(1)$ | $[0]_x$ | | 3.22 | |
| $Cl(2) \cdots H(1)$ | $\begin{bmatrix} 0 \end{bmatrix}_x$ | | 2.92 | |
| $Cl(2) \cdots H(3)$ | $[1]_x$ | | 2.96 | |
| $H(5) \cdots Cl(2)$ | $[1]_x$ | | 2.98 | |
| $Cl(1) \cdots O(1)$ | $[2]_{x}$ | | 3.60 | |
| $Cl(2) \cdots O(2)$ | $[2]_{y}$ | | 3.43 | |
| $Cl(2) \cdots Cl(2)$ | [3] | | 3.30 | |
| | | | | Angle |
| $C(1)-Cl(1)\cdots Cl$ | (2) | [3] | | 147.3 |

* The four molecules are numbered [0], [1], [2] and [3], with:

| [0] | at | <i>x</i> , | y, | z; |
|-----|---------------|---------------------|---------------------|--------------------|
| [1] | \mathbf{at} | $\frac{1}{2} - x$, | -y, | $\frac{1}{2} + z;$ |
| [2] | \mathbf{at} | $\frac{1}{2} + x$, | $\frac{1}{2} - y$, | -z; |
| [3] | \mathbf{at} | -x, | $\frac{1}{2} + y$, | $\frac{1}{2}-z$. |

Translation of the unit cell along one of the crystallographic axes is indicated by a subscript.

The main short approach distances occur between molecules related by the unit translation along [100], indicating a close stacking in this direction. Of particular interest are the separations of $C(3) \cdots O(2)$ $(3 \cdot 22 \text{ Å})$ and of $O(2) \cdots C(1)$ $(3 \cdot 14 \text{ Å})$. Similar or shorter non-hydrogen bonded $C \cdots O$ distances have been reported in recent years (Sime & Abrahams, 1960; Chu, Jeffrey & Sakurai, 1962; Romers, 1963); however, all of these were found to occur between molecules in which at least one of the atoms involved takes part in a π -bonded system.

The packing is less dense in all other directions, in accord with the easy cleavage of the crystals along planes in the [100] zone. Noteworthy is the short distance $(3\cdot303\pm0\cdot007\text{ Å})$ that occurs between axial and equatorial chlorine atoms approximately in the directions [011], [011] etc. The approach is even closer than that found in trans-2,5-dichloro-1,4-dioxane $(3\cdot38\text{ Å})$ and in β -2-chloro-p-benzoquinone-4-oxime acetate $(3\cdot33\text{ Å}, \text{Fischman}, \text{MacGillavry \& Romers},$ 1961). A short Cl \cdots Cl distance $(3\cdot34\pm0\cdot04\text{ Å})$ also occurs in solid chlorine (Collin, 1952).



Fig. 4. Crystal structure of cis-2,3-dichloro-1,4-dioxane. The four symmetry-related molecules are numbered [0], [1], [2], and [3] (see Table 6). (a) Projection of the structure along [100], molecules shown in dotted outlines belong to a layer of unit cells translated -1 along x. (b) Projection along [010]. For clarity, molecule [2]-x is not shown.

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Proton magnetic resonance

It is of interest to note that a single crystal of cis-2,3-dichloro-1,4-dioxane contains either right-handed or left-handed molecules. Here we have an example of an asymmetric conformational transformation induced by crystallization (cf. Havinga, 1954). The solution is optically inactive, because at room temperature in the liquid phase the equilibrium between the two enantiomeric conformations (ea and ae) is readily established; *i.e.* an axial atom becomes an equatorial one and vice versa by chair-chair interconversion.

Proton magnetic resonance (p.m.r.) spectra were taken at various temperatures in order to verify these conclusions (Smidt & Altona, 1962). At temperatures between +20 and -100 °C only one sharp resonance peak is observed for the two non-equivalent protons at C(1) and C(2); hence the ring inversion is rapid to the extent that the chemical shifts of equatorial H and axial H are averaged to a single value. This 'collapsed' signal is found at higher field $(\tau=4.37)$ than that of the (ee) protons in the trans-2,3-isomer ($\tau = 4.10$) in accordance with a well known rule (Lemieux, Kullnig & Moir, 1958). At -118 °C the inversion rate is reduced sufficiently to effect an appreciable broadening of the resonance signal and at -128 °C two lines are observed. This temperature is surprisingly low. Recent p.m.r. studies have shown that splitting of the appropriate resonance peaks of cyclohexane (Jensen, Noyce, Sederholm & Berlin, 1960; Moniz & Dixon, 1961) and of several halogenocyclohexanes (Berlin & Jensen, 1960; Reeves & Strømme, 1960, 1961; van Dort & Sekuur, 1963) takes place at temperatures between -60 and -80 °C.

After completion of this work a publication appeared (Caspi, Wittstruck & Piatak, 1962) in which a conformation of *cis*-2,3-dichloro-1,4-dioxane in solution was derived from p.m.r. data. The spectrum is in excellent agreement with that taken at 20 °C by Smidt & Altona (1962). However, the conclusion that the molecule probably exists in solution as a rigid boat conformation is unjustified as it was not based upon measurements performed at a sufficiently low temperature. Furthermore, the infrared spectra of the *cis* compound in the solid state (KBr disc) and in solution (CS₂) are practically identical (Altona, 1963). This fact also indicates that the (*ea*) chair form occurring in the crystal is also the preferred one in solution.

We are grateful to Prof. E. Havinga for suggesting this research and for his kind interest. We wish to thank Dr Carolyn Knobler for performing many tedious calculations during the early stages of the analysis, Miss E. Alzerda and Mr E. van Heijkoop for technical assistance, and Dr A. Vos and Mr H. Schurer (University of Groningen, The Netherlands) for their valuable cooperation in carrying out the calculations on the ZEBRA computer.

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